

Clebsch-Gordan Coefficient Example Phys 402

The eigenfunctions of J^2 can be expressed as linear combinations of states with different values of m_ℓ and m_s using the world-famous Clebsch-Gordan coefficients

$(C_{m_\ell m_s m_j}^{\ell s j})$ as:

$$|j m_j\rangle = \sum_{m_\ell+m_s=m_j} C_{m_\ell m_s m_j}^{\ell s j} |\ell m_\ell\rangle |s m_s\rangle \quad [4.185]$$

Where the ket $|\ell m_\ell\rangle$ represents the spherical harmonics $Y_\ell^{m_\ell}$. The C-G coefficient values are given in Table 4.8 on page 179 of Griffiths. Remember that all of the coefficients should appear under a square root, with the minus sign (if any) out front.

Now for an example of how to construct states that are simultaneous eigenfunctions of L^2 , S^2 , J^2 and J_z . Take the case again of hydrogen with $\ell = 1$ and spin $s = 1/2$. How do we find the state with $j = 3/2$ and $m_j = -1/2$ in terms of the $Y_\ell^{m_\ell}$ and spinors? Look at the $1 \times 1/2$ CG Table on page 179. We are led to this table because we are combining an angular momentum vector with $\ell = 1$ and spin vector with $s = 1/2$.

$1 \times 1/2$	$3/2$				
	$+3/2$	$3/2$	$1/2$		
	$+1$	$+1/2$	1	$+1/2$	$+1/2$
		$+1$	$-1/2$	$1/3$	$2/3$
		0	$+1/2$	$2/3$	$-1/3$
			0	$-1/2$	$2/3$
			-1	$+1/2$	$1/3$
				-1	$-1/2$
					1
2×1	3	3	2		
	$+3$	3	2		

Now look under the column labeled “ $\begin{matrix} 3/2 \\ -1/2 \end{matrix}$ “. It says:

$$\left| \begin{matrix} 3 \\ 2 \\ -1/2 \end{matrix} \right\rangle = \sum_{m_\ell+m_s=-1/2} C_{m_\ell m_s -1/2}^{1 \ 1/2 \ 3/2} \left| 1 m_\ell \right\rangle \left| \frac{1}{2} m_s \right\rangle$$

$$\left| \begin{matrix} 3 \\ 2 \\ -1/2 \end{matrix} \right\rangle = \sqrt{\frac{2}{3}} \left| 1 \ 0 \right\rangle \left| \frac{1}{2} \ -\frac{1}{2} \right\rangle + \sqrt{\frac{1}{3}} \left| 1 \ -1 \right\rangle \left| \frac{1}{2} \ \frac{1}{2} \right\rangle$$

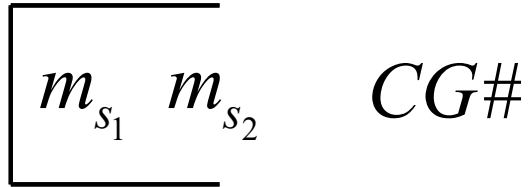
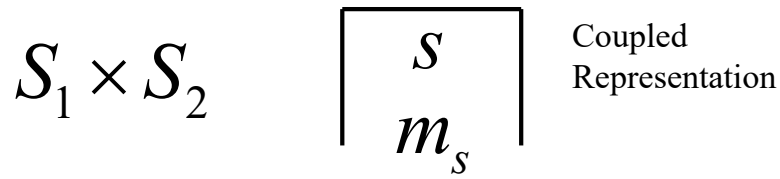
This can be written in a more familiar way in terms of spherical harmonics and spinors as:

$$\langle \vec{r} | \begin{matrix} 3 \\ 2 \\ -1/2 \end{matrix} \rangle = \sqrt{\frac{2}{3}} Y_1^0(\theta, \phi) \chi_- + \sqrt{\frac{1}{3}} Y_1^{-1}(\theta, \phi) \chi_+$$

One can move back and forth between the coupled and un-coupled representations using the Clebsch-Gordan table on page 179. Here is the schematic layout for the CG table for combining two spins (called \vec{S}_1, \vec{S}_2) to form a total spin $\vec{S} = \vec{S}_1 + \vec{S}_2$ (S^2 has eigenvalue $s(s+1)\hbar^2$):

General Schematic of the C-G Table

$$\vec{S} = \vec{S}_1 + \vec{S}_2$$



Un-Coupled Representation